Dynamic Modeling and Analysis of Vertical Vibration Reduction System for Passenger Train

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1. Introduction

A gap in the mechanical system will make the system produce collision and impact, which will affect the performance and safe operation of the mechanical system. For example, the impact between the wheels and the track in the running of high-speed trains intensifies the vibration of the trains and affects the running stability and comfort of the trains. Therefore, the study of vibro-impact is of great significance to reduce the collision, impact and abrasion of mechanical system, as the same time to improve the safety, lifetime and efficiency of the mechanical system. In recent years, the theory and application of vibro-impact system and gap system have made rapid progress [1-3].

The existing vehicle vertical vibration reduction system is mainly designed according to the track (road surface) irregularity [4, 5]. Wei et al. [6] used the proposed dynamic model for a safety analysis and a vibration-reduction evaluation to theoretically validate the feasibility of semi-active magneto-rheological steel-spring FST. Cai et al. [7] studied the vibration control effect of long elastic sleeper track in subways through experiment. However, due to the existence of the wheel and rail gap, even if the train is running on completely smooth track, Chaos vibration will occur. There are many studies on wheel-rail impact and vehicle control [8-12]. Choi et al. [13] predicted the vibration of trains running on ballasted track by measuring the impact factor of track. However, these papers rarely analyze train vibration from the perspective of vibro-impact. In recent years, there are much theoretical research on vibro-impact [14-18]. Liu et al. [19] studied the dynamics of a capsule system in various friction environments. Yue et al. [20] focused on the coexistence of strange non-chaotic attractors (SNAs) and a novel mixed attractor in a periodically driven three-degree-of-freedom vibro-impact system with symmetry. However, few of them are combined with engineering practice. Therefore, in this paper, based on the wheel-rail impact vibration, the stiffness and natural damping of the carriage are considered. The dynamic model of the vertical vibration damping system of the train is built. The influence of different damping system parameters on the vibration of the running train is studied.

$$m_1\ddot{x}_1 + 2\zeta_1\dot{x}_1 + k_1x_1 = f_{a0}\sin(\omega t + \tau).$$

2. System model

Assuming that the passenger train’s wheel-rail impact is a rigid impact and the motion in the vertical direction is considered only, the model of a single wheel vibro-impact system is shown in Fig. 1. The meanings of the symbols in Fig. 1 are shown in Table 1. Natural damping and carriage stiffness are connected to imaginary inertial spaces. When the displacement of the wheel minus the displacement of the rail is equal to the gap, wheel and rail impact. After impacting, wheel and rail get new speed and then impact again and again.

![Dynamic model of a single wheel vibro-impact system for passenger train](image)

As shown in Fig. 1, the dimensionless differential equation where the system does not impact with each other can be written as:

$$\begin{bmatrix}
1 & 0 & 0 & \dot{x}_1 \\
0 & m_1 & 0 & \ddot{x}_1 \\
0 & 0 & m_1 & \ddot{x}_1 \\
\end{bmatrix}
+ \begin{bmatrix}
2\zeta_1 & -2\zeta_1 & 0 \\
-2\zeta_2 & 2\zeta_2 & -2\zeta_2 \\
0 & 2\zeta_2 & 2\zeta_2 \\
\end{bmatrix}
\begin{bmatrix}
x_1 \\
\dot{x}_1 \\
\dot{x}_1 \\
\end{bmatrix}
+ \begin{bmatrix}
k_1 & -k_1 & 0 \\
-k_2 & k_1+k_2 & -k_2 \\
0 & -k_2 & k_2+k_1 \\
\end{bmatrix}
\begin{bmatrix}
x_1 \\
\dot{x}_1 \\
\dot{x}_1 \\
\end{bmatrix}
= \begin{bmatrix}
f_{a0} \\
\end{bmatrix}
\sin(\omega t + \tau),$$

(1)

In the formula, dimensionless quantities are:
When the gap is 0, the motion equation of the wheel and rail is:

\[ \ddot{x}_t + m_1 \dot{x}_t = \dot{x}_i + m_4 \dot{x}_4, \]

(4)

\[ R = \frac{\dot{x}_4 - \dot{x}_i}{\dot{x}_t - \dot{x}_4}. \]

(5)

The dimensionless instantaneous velocity before and after the impact of the wheel can be represented by \( \dot{x}_i \) and \( \dot{x}_t \) respectively. The dimensionless instantaneous velocity before and after the rail impact can be expressed by \( \dot{x}_4 \) and \( \dot{x}_t \) respectively. \( R \) is the coefficient of restitution. Eqs (4) and (5) can be obtained:

\[ \ddot{x}_i = \frac{1 - R m_4}{1 + m_4} \dot{x}_i + m_1 (1 + R) \dot{x}_4, \]

(6)

\[ \ddot{x}_t = \frac{1 + R}{1 + m_4} \dot{x}_4 - \frac{m_i - R}{1 + m_4} \dot{x}_i. \]

(7)

The semi-analytical solution of the system is then obtained through a series of derivation. Select Poincaré section: \( \theta = \omega t \), and the Poincaré map of periodic motion of the system can be established by selecting the section:

\[ \Delta X' = f(\mathbf{v}, \Delta X), \]

(8)

where: \( \Delta X' = (x_{10}, \dot{x}_i, x_{30}, \dot{x}_3, x_{50}, \dot{x}_5, x_{70}, \dot{x}_7)^{T}, \mathbf{v} \in \mathbb{R}^4 \), are fixed points in the periodic collision motion process of the system on the Poincaré section \( \sigma \), \( \Delta X \) and \( \Delta X' \) are the relevant disturbance quantity of the fixed point on the section \( \sigma \).

When appropriate system parameters are selected, the symbol \( q = pn \) is usually used to represent the periodic motion of the vibro-impact system (\( q = pn \) does not represent the rational number in general, but only the symbol here), \( n \) is used to represent the number of periods of force, and \( p \) is used to represent the number of impacts. The periodic motion of \( q = 1/n \) refers to: directly set the instantaneous time \( t \) after the impact of the two oscillators \( M_1 \) and \( M_4 \) to be 0. It is not difficult to find that the dimensionless time \( t \) is just equal to \( 2\pi/\omega(n = 1, 2, \ldots) \) at the moment when the two oscillators \( M_1 \) and \( M_4 \) impact with each other. By shifting the origin of \( \theta \) coordinate to a collision point, the boundary conditions of periodic motion of the system can be known:

\[
\begin{align*}
x_{10} = x_i(0) = x_i(2\pi/\omega), x_i(0) - x_i(0) &= b \\
\dot{x}_{10} = \dot{x}_i(0) = 1 - R m_4 \dot{x}_i(2\pi/\omega) + m_1 (1 + R) \dot{x}_i(2\pi/\omega) \\
\dot{x}_{30} = \dot{x}_3(0) = x_3(2\pi/\omega), \dot{x}_{20} = \dot{x}_2(0) &= \dot{x}_2(2\pi/\omega) \\
\dot{x}_{50} = \dot{x}_5(0) = 1 + R \dot{x}_5(2\pi/\omega) + m_4 - R \dot{x}_5(2\pi/\omega) \\
\dot{x}_{70} = \dot{x}_7(0) &= \dot{x}_7(2\pi/\omega) \\
x_{10} = x_i(0) = x_i(2\pi/\omega), \dot{x}_{10} = \dot{x}_i(0) &= \dot{x}_i(2\pi/\omega)
\end{align*}
\]

(9)

Thus, the system parameters of the vibro-impact system moving in \( q = 1/n \) period can be obtained, where \( \dot{x}_i(2\pi/\omega) = \dot{x}_i \) and \( \dot{x}_i(2\pi/\omega) = \dot{x}_i \) are the expressions of the impact velocity of the system. The linearized matrix (formal solution of Jacobian matrix) of Poincaré mapping at a fixed point can be written as:

\[
Df(\mathbf{v}, X') = \frac{\partial f(\mathbf{v}, \Delta X)}{\partial \Delta X} \bigg|_{(\mathbf{v}, X')}.
\]

(10)

### 3. Dynamics simulation and analysis of the system

According to the actual train parameters (\( K_1 \) is between 0.2 and 1.5 MN/m, \( K_2 \) is between 0.1 and 0.6 MN/m, \( C_1 \) is between 10 and 30 kN/s, \( C_2 \) is between 15 and 45 kN/s), the simulation parameters are selected as shown in the Table 2. The first column of the table is the variable symbols, the second column is the first set of parameter values, and the third column is the second set of parameter values. The second set of parameters differs from the first set in value \( \zeta_1 \) only.

<table>
<thead>
<tr>
<th>Table 2</th>
<th>Simulation parameters</th>
</tr>
</thead>
<tbody>
<tr>
<td>( m_2, m_3, m_4 )</td>
<td>1.7 1.7 0.6</td>
</tr>
<tr>
<td>( k_2, k_3, k_4 )</td>
<td>0.8 1.2 1.5</td>
</tr>
<tr>
<td>( \zeta_2, \zeta_3, \zeta_4 )</td>
<td>1.2 1.5 0.5</td>
</tr>
<tr>
<td>( \zeta_1 R \cdot b )</td>
<td>0.6 0.7 1.5</td>
</tr>
</tbody>
</table>

Fig. 2 shows the Poincaré map generated after the
simulation of the first set of parameters. Fig. 2a shows a single point on the Poincaré map. Fig. 2b shows two points on the Poincaré map. Fig. 3c shows four points on the Poincaré map, and Fig. 2d shows an infinite number of points on the Poincaré map.

Fig. 2 Poincaré maps of the first set of parameters

![Fig. 2 Poincaré maps of the first set of parameters](image)

Fig. 3 shows the acceleration time diagram formed after the simulation of the first parameter. As shown in Fig. 3, a, the acceleration presents a single-period characteristic, with its maximum value less than 0.4. As shown in Fig. 3, b, the acceleration presents a double-periodic characteristic with an amplitude of close to 0.4. As shown in Fig. 3, c, the acceleration presents a 4-period characteristic, whose amplitude is greater than 0.4. As shown in Fig. 3, d, the acceleration curve has no periodicity, and its amplitude is close to 0.5.

Choosing the first set of parameters to the simulation, when \( \omega = 2.23 \), the system has a stable periodic motion \( q = 1/1 \), as shown in Fig. 2, a and Fig. 3, a. As \( \omega \) increases, the system doubles the bifurcation as \( q = 2/2 \) periodic motion, as shown in Fig. 2, b and Fig. 3, b. \( \omega \) keeps going up to the \( \omega = 2.249 \), the system doubles to periodic motion \( q = 4/4 \), as shown in Fig. 2, c and Fig. 3, c. Finally, as \( \omega \) continues to increase, the system goes into chaos motion, as shown in Fig. 2, d and Fig. 3, d.

The simulation results of the second set of parameters are shown in Fig. 4. Fig. 4, a and Fig. 4, c are Poincaré maps, Fig. 4, b and Fig. 4, d are acceleration time diagrams.

When the second set of parameters is selected for simulation, compared with the first set of parameters, only the value \( \zeta_1 \) of the system is reduced. The reduction of \( \zeta_1 \) also means that the damping of the primary suspension is reduced. By comparing Fig. 2, a with Fig. 4, a, it is found that the motion of the system changes from stable period 1 to period 4 when the value of \( \zeta_1 \) decreases with the constant \( \omega \). By comparing Fig. 3, a and Fig. 4, b, it is also not difficult to find that the amplitude of wheel’s acceleration increases with the decrease of \( \zeta_1 \).

By comparing Fig. 2, b with Fig. 4, c, it is found that the system doubles from 2 period to 4 period with the decrease of \( \zeta_1 \). Similarly, the amplitude of wheel’s acceleration is also increased, as shown in Fig. 3, b and Fig 4, d.

4. Engineering verification and discussion

During the operation of the train, the mass of the whole train will change due to the change in the number of passengers. Track stiffness and damping will change due to the change of road conditions. As the temperature changes, the damping and stiffness of the carriage will change. The initial sensitivity of a chaotic system can cause a completely different vibration effect. Once, an abnormal vibration occurred in one carriage of the train, and the periodic amplitude of vertical acceleration of wheelset was found to increase during detection (The test point was placed on the wheelset bearings), as shown in Fig. 5. We thought there may have been structural damage to the wheels or the bearing and other bogie components, but no structural damage was found when the further
Fig. 3 Acceleration time diagrams of the first set of parameters

Fig. 4 Simulation results of the second set of parameters
inspection was conducted, so we suspected that the system had bifurcation. The primary suspension shock absorber of this car was exchanged with a new shock absorber. When the train was running again at the same speed of 60 km/h, no abnormal vibration occurred in this car. Fig. 6 shows the vertical acceleration of wheelset after the shock absorber was exchanged. Later we checked the old shock absorber and found that there was a slow leakage of oil which reduced the damping. This engineering practice verified the correctness of our simulation.

![Graph](image)

Fig. 5 Vertical vibration acceleration of wheelset of abnormal vibration carriage

![Graph](image)

Fig. 6 Vertical vibration acceleration of wheelset after shock absorber replacement

5. Conclusions

This paper analyzes the vertical vibration damping system of the passenger train. The dynamic model of the passenger train’s vertical vibration damping system is built as a four-degree-of-freedom vibro-impact system model. The simulation shows that the train may enter into a chaotic state in the actual operation process. The engineering practice shows that the vertical vibration of the train will lose its stability with the decrease of primary suspension damping, and the simulation results are also verified. It has been proved that the vibration reduction design of the train should consider not only the track unevenness, but also the multi-period vibration, even chaos vibration caused by wheel-rail vibro-impact.

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References


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